HOW LIGHT IS DARK?

V.J. Stenger

June, 1987

It has been proposed that photomultipler tubes in the sea are a very sensitive way to search for the class of costs which have nuclear candidates in the mass range $10^{-10} < M < 10$ g which have nuclear density and are no smaller than atoms. Here some details are given on the light production mechanism and the luminous efficiency of seawater is calculated.

Following De Rujula and Glashow, 2 it is assumed that the particle emits a cylindrical shock wave as it travels through the water with a speed v. The rate of expansion of the shock is given by

$$R(t) = (2nT/m)^{1/2}$$
 (1)

where m is the molecular mass, n is the number of degrees of freedom (taken as 3), and the temperature T of the shock at the time t is given by

$$T(t) = T_{O}(t_{O}/t) \tag{2}$$

where $T_0 = mv^2/n$ is the intitial temperature (natural units are used). The radius of the shock at a time t is then given by

$$R(t) = R_0(t/t_0)^{1/2}$$
 (3)

where R_0 is the radius of the particle and $t_0 = R_0/\sqrt{8}v$ is an integration constant which specifies that time at which the shock starts.

Light is assumed to be emitted with the black-body spectrum,

$$dp \qquad \omega^3 \qquad 1$$

$$d\omega da \qquad (2\pi)^2 e^{\omega/T} - 1$$
(4)

and the optical energy emitted in the band $(\omega_{\min}, \omega_{\max})$ is

$$dE_{y} = \omega_{\text{max}} \infty \qquad dp$$

$$--- = \int d\omega \int dt \ 2\pi \ R(t) \ ---$$

$$dx = \omega_{\text{min}} \quad t_{0} \qquad (5)$$

In all the cases here, the radius R of the shock exceeds the mean free path of the particle in the water.

The luminous efficiency then is defined as

$$\eta = (dE_{\gamma}/dx)/(dE/dx)$$
 (6)

where dE/dx = $\rho v^2 A$, $A = \pi R_0^2$.

De Rujula and Glashow assume that the medium is highly

transparent and integrate (5) analytically, assuming also that $T >> \omega_{max}$. Since seawater is not transparent, and in fact has a sharp spectral acceptance, this may not be a good assumption for our purposes. Thus I have integrated (5) numerically. I find that η depends on both the assumed spectral band $(\lambda_{min}, \lambda_{max})$ and the dark particle mass M, with a peak at $M = 5 \times 10^{-11}$ g. The results are shown in Fig. 1. For the whole visible band (350,650) nm, I obtain $\eta = 2 \times 10^{-5}$ at the peak, in agreement with De Rujula and Glashow. To conservatively approximate the response of seawater I assume a smaller band, (400,500) nm, and this is also shown in Fig. 1.

Since the attenuation of light in water depends sharply on wavelength, a more precise calculation would compute a luminous efficiency as a function of ω and propagate each ω separately to the array. However, this should be an adequate approximation for most purposes..

references

- 1. V.J. Stenger, HDC-3-87
- 2. A. De Rujula and S.L. Glashow, Nature 312, 734 (1984).

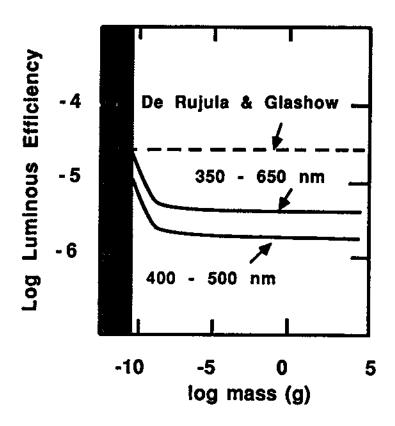


Fig. 1. The luminous efficiency as a function of dark particle mass, for two spectral bands. The 400 - 500 nm band was used in this report. The calculation of De Rujula and Glashow gives $2x10^{-5}$ independent of mass, for highly transparent water.